PHASELESS INVERSE SCATTERING
IN THE ONE-DIMENSIONAL CASE

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Abstract We consider the one-dimensional Schrödinger equation with a potential satisfying the standard assumptions of the inverse scattering theory and supported on the half-line $x \geq 0$. For this equation at fixed positive energy we give explicit formulas for finding the full complex valued reflection coefficient to the left from appropriate phaseless scattering data measured on the left, i.e. for $x < 0$. Using these formulas and known inverse scattering results we obtain global uniqueness and reconstruction results for phaseless inverse scattering in dimension $d = 1$.

1 Introduction

We consider the one-dimensional Schrödinger equation

$$-\frac{d^2}{dx^2}\psi + v(x)\psi = E\psi, \quad x \in \mathbb{R}, \quad E > 0, \tag{1.1}$$

where v satisfies the standard assumptions of the inverse scattering theory (see [5]) and is supported on the half-line $x \geq 0$. More precisely, we assume that

$$v \text{ is real} - \text{valued}, \quad v \in L_1^1(\mathbb{R}), \tag{1.2}$$

$$v(x) \equiv 0 \text{ for } x < 0,$$

where

$$L_1^1(\mathbb{R}) = \{u \in L^1(\mathbb{R}) : \int_{\mathbb{R}} (1 + |x|)|u(x)|dx < \infty\}. \tag{1.3}$$

For equation (1.1) we consider the scattering solution $\psi^+ = \psi^+(\cdot, k)$, $k = \sqrt{E} > 0$, continuous and bounded on $\mathbb{R}$ and specified by the following asymptotics:

$$\psi^+(x, k) = \begin{cases} e^{ikx} + s_{21}(k)e^{-ikx} & \text{as } x \to -\infty, \\ s_{22}(k)e^{ikx} + o(1) & \text{as } x \to +\infty, \end{cases} \tag{1.4}$$

for some a priori unknown s_{21} and s_{22}. In addition, the coefficients s_{21} and s_{22} arising in (1.4) are the reflection coefficient to the left and transmission coefficient to the right, respectively, for equation (1.1).

In order to find ψ^+ and s_{21}, s_{22} from v one can use well-known results of the one-dimensional direct scattering theory, see e.g. [5]. And properties of ψ^+, s_{21}, s_{22} are known in detail, see [4], [5], [7], [12].

In particular, it is well-known that

$$|s_{21}(k)|^2 + |s_{22}(k)|^2 = 1, \quad |s_{21}(k)|^2 < 1, \quad k > 0. \tag{1.5}$$

Let

$$\mathbb{R}_+ =]0, +\infty[, \quad \mathbb{R}_- =]-\infty, 0[. \tag{1.6}$$

We consider the following two types of scattering data measured on the left for equation (1.1): (a) $s_{21}(k)$ and (b) $\psi^+(x,k)$, $x \in X_- \subseteq \mathbb{R}_-$, where $k = \sqrt{E} > 0$.

In addition, we consider the following inverse scattering problems:

Problem 1.1a. Reconstruct potential v on $\mathbb{R}$ from its reflection coefficient s_{21} on $\mathbb{R}_+$.

Problem 1.1b. Reconstruct potential v on $\mathbb{R}$ from its scattering data $\psi^+(x, \cdot)$ on $\mathbb{R}_+$ at fixed $x \in \mathbb{R}_-$.

Problem 1.2a. Reconstruct potential v on $\mathbb{R}$ from its phaseless scattering data $|s_{21}|^2$ on $\mathbb{R}_+$.

Problem 1.2b. Reconstruct potential v on $\mathbb{R}$ from its phaseless scattering data $|\psi^+|^2$ on $X_- \times \mathbb{R}_+$ for some appropriate X_-.

Problem 1.2c. Reconstruct potential v on $\mathbb{R}$ from its phaseless scattering data $|s_{21}|^2$ on $\mathbb{R}_+$ and $|\psi^+|^2$ on $X_- \times \mathbb{R}_+$ for some appropriate X_-.

Note that in quantum mechanical scattering experiments in the framework of model described by equation (1.1) the phaseless scattering data $|s_{21}|^2$, $|\psi^+|^2$ of Problems 1.2a-1.2c can be measured directly, whereas the complex-valued scattering data s_{21}, ψ^+ of Problems 1.1a, 1.1b are not accessible for direct measurements. Therefore, Problems 1.2 are of particular applied interest in the framework of inverse scattering of quantum mechanics. However, Problems 1.1 are much more considered in the literature than Problems 1.2. See [3]-[5], [12]-[14], [17] and references therein in connection with Problem 1.1a and [1], [11] in connection with Problem 1.2a and its modification.

In particular, work [17] gives global uniqueness and reconstruction results for Problem 1.1a; see also [2], [6] and references given in [2]. And, obviously, Problem 1.1b is reduced to Problem 1.1a by the formula

$$s_{21}(k) = e^{ikx}\psi^+(x,k) - e^{2ikx}, \quad x \in \mathbb{R}_-, \quad k \in \mathbb{R}_+. \tag{1.7}$$

On the other hand, for Problem 1.2a it is well known that the phaseless scattering data $|s_{21}|^2$ on $\mathbb{R}_+$ do not determine v uniquely, in general. In particular, we have that

$$s_{21,y}(k) = e^{2iky}s_{21}(k), \tag{1.8}$$

$$|s_{21,y}(k)|^2 = |s_{21}(k)|^2, \quad k \in \mathbb{R}_+, \quad y \in \mathbb{R},$$

where s_{21} is the reflection coefficient to the left for v and $s_{21,y}$ is the reflection coefficient to the left for v_y, where

$$v_y(x) = v(x - y), \quad x \in \mathbb{R}, \quad y \in \mathbb{R}. \tag{1.9}$$

This well-known nonuniqueness result on Problem 1.2a motivates our considerations of Problems 1.2b and 1.2c.

In the present work we continue studies of [16]. We recall that article [16] gives, in particular, explicit asymptotic formulas for finding the complex-valued scattering amplitude from appropriate phaseless scattering data for the Schrödinger equation at fixed energy in dimension $d \geq 2$. Then using these formulas related phaseless inverse scattering problems are reduced in [16] in dimension $d \geq 2$ to well developed inverse scattering from the complex-valued scattering amplitude. And, as a corollary, [16] contains different results on inverse scattering without phase information in dimension $d \geq 2$.

In connection with recent results on phaseless inverse scattering in dimension $d \geq 2$, see also [8]-[10], [15] and references therein.

In order to present results of the present work we use the notations:

$$S_1(x_1, x_2, k) = \{|s_{21}(k)|^2, |\psi^+(\cdot, k)|^2 \quad \text{on} \quad X_-\}, \tag{1.10}$$

$$\text{where} \quad X_- = \{x_1, x_2 \in \mathbb{R}_- : \quad x_1 \neq x_2\}, \quad k \in \mathbb{R}_+;$$

$$S_2(x_1, x_2, x_3, k) = |\psi^+(\cdot, k)|^2 \quad \text{on} \quad X_-, \tag{1.11}$$

$$\text{where} \quad X_- = \{x_1, x_2, x_3 \in \mathbb{R}_- : \quad x_i \neq x_j \text{ if } i \neq j\}, \quad k \in \mathbb{R}_+;$$

$$S_3(x, k) = \{|\psi^+(x, k)|^2, \frac{d|\psi^+(x, k)|^2}{dx}\}, x \in \mathbb{R}_-, \ k \in \mathbb{R}_+. \tag{1.12}$$

Using these notations the main results of the present work can be summarized as follows:

(A1) We give explicit formulas for finding $s_{21}(k)$ from $S_1(x_1, x_2, k)$ for fixed x_1, x_2 and k, where $x_1 \neq x_2 \ mod(\pi(2k)^{-1})$; see Theorem 2.1 of Section 2.

(A2) We give explicit formulas for finding $s_{21}(k)$ from $S_2(x_1, x_2, x_3, k)$ for fixed x_1, x_2, x_3 and k, where $x_i \neq x_j \ mod(\pi k^{-1})$ if $i \neq j$; see Theorem 2.2 of Section 2.

(A3) We give explicit formulas for finding $s_{21}(k)$ from $S_3(x, k)$ for fixed x and k; see Theorem 2.3 of Section 2.

(B) We give global uniqueness and reconstruction results (1) for finding v on $\mathbb{R}$ from $S_1(x_1, x_2, \cdot)$ on $\mathbb{R}_+$ for fixed x_1, x_2, (2) for finding v on $\mathbb{R}$ from $S_2(x_1, x_2, x_3, \cdot)$ on $\mathbb{R}_+$ for fixed x_1, x_2, x_3, and (3) for finding v on $\mathbb{R}$ from $S_3(x, \cdot)$ on $\mathbb{R}_+$ at fixed x; see Theorem 2.4 of Section 2.

Note that results of (B1)-(B3) follow from (A1)-(A3) and the aforementioned results of [17] on Problem 1.1a. In addition, the results of (B1) are global results on Problem 1.2c and the results of (B2), (B3) are global results on Problem 1.2b.

The main results of the present work are presented in detail in Section 2.

2   Main results

We represent s_{21} of (1.4) as follows:

$$s_{21}(k) = |s_{21}(k)|e^{i\alpha(k)}, \quad k \in \mathbb{R}_+. \tag{2.1}$$

We consider

$$a(x, k) = |\psi^+(x, k)|^2 - 1, \quad x \in \mathbb{R}_-, \quad k \in \mathbb{R}_+, \tag{2.2}$$

where ψ^+ is the scattering solutions of (1.4).

Theorem 2.1. *Let potential v satisfy (1.2) and s_{21}, a be the functions of (1.4), (2.2). Let $x_1, x_2 \in \mathbb{R}_-$, $k \in \mathbb{R}_+$, $x_1 \neq x_2 \mod (\pi(2k)^{-1})$. Then:*

$$|s_{21}| \begin{pmatrix} \cos \alpha \\ \sin \alpha \end{pmatrix} = (2\sin(2k(x_2 - x_1)))^{-1} \times$$

$$\begin{pmatrix} \sin(2kx_2) & -\sin(2kx_1) \\ -\cos(2kx_2) & \cos(2kx_1) \end{pmatrix} \begin{pmatrix} a(x_1, k) - |s_{21}|^2 \\ a(x_2, k) - |s_{21}|^2 \end{pmatrix}, \tag{2.3}$$

where $|s_{21}| = |s_{21}(k)|$, $\alpha = \alpha(k)$.

Theorem 2.1 is proved in Section 3.

One can see that Theorem 2.1 gives explicit formulas for finding $s_{21}(k)$ from $S_1(x_1, x_2, k)$ of (1.10) for fixed x_1, x_2 and k, where $x_1 \neq x_2 \mod(\pi(2k)^{-1})$.

Theorem 2.2. *Let potential v satisfy (1.2) and s_{21}, a be the functions of (1.4), (2.2). Let $x_1, x_2, x_3 \in \mathbb{R}_-$, $k \in \mathbb{R}_+$, $x_i \neq x_j \mod(\pi k^{-1})$ if $i \neq j$. Then:*

$$|s_{21}| \begin{pmatrix} \cos \alpha \\ \sin \alpha \end{pmatrix} = (8(\sin(k(x_2 - x_3))\sin(k(x_2 - x_1))\sin(k(x_1 - x_3)))^{-1} \times \tag{2.4}$$

$$\begin{pmatrix} \sin(2kx_3) - \sin(2kx_1) & -\sin(2kx_2) + \sin(2kx_1) \\ -\cos(2kx_3) + \cos(2kx_1) & \cos(2kx_2) - \cos(2kx_1) \end{pmatrix} \begin{pmatrix} a(x_2, k) - a(x_1, k) \\ a(x_3, k) - a(x_1, k), \end{pmatrix}$$

where $|s_{21}| = |s_{21}(k)|$, $\alpha = \alpha(k)$.

Theorem 2.2 is proved in Section 3.

One can see that Theorem 2.2 gives explicit formulas for finding $s_{21}(k)$ from $S_2(x_1, x_2, x_3, k)$ of (1.11) for fixed x_1, x_2, x_3 and k, where $x_1 \neq x_2 \mod(\pi k^{-1})$ if $i \neq j$.

Theorem 2.3. *Let potential v satisfy (1.2) and s_{21}, ψ^+ be the functions of (1.4). Then the following formulas hold:*

$$Re\,(s_{21}(k)e^{-ikx}) = -1 + (|\psi^+(x, k)|^2 - |Im\,(s_{21}(k)e^{-ikx})|^2)^{1/2}, \tag{2.5}$$

$$Im\,(s_{21}(k)e^{-ikx}) = \frac{1}{4k}\frac{d|\psi^+(x, k)|^2}{dx}, \tag{2.6}$$

where $x \in \mathbb{R}_-$, $k \in \mathbb{R}_+$, and $(\cdot)^{1/2} > 0$ in (2.5).

Theorem 2.3 is proved in Section 3.

One can see that Theorem 2.3 gives explicit formulas for finding $s_{21}(k)$ from $S_3(x, k)$ of (1.12).

As corollaries of Theorems 2.1, 2.2, 2.3 and results of [17], we obtain the following global uniqueness and reconstruction results on phaseless inverse scattering for equation (1.1):

Theorem 2.4. *Let potential v satisfy (1.2) and S_1, S_2, S_3 be the phaseless scattering data of (1.10), (1.11), (1.12). Then: (1) $S_1(x_1, x_2, \cdot)$ on $\mathbb{R}_+$, for fixed x_1, x_2, uniquely determine v on $\mathbb{R}$ via formulas (2.1)-(2.3) and results of [17]; (2) $S_2(x_1, x_2, x_3, \cdot)$ on $\mathbb{R}_+$, for fixed x_1, x_2, x_3, uniquely determine v on $\mathbb{R}$ via formulas (2.1), (2.2), (2.4) and results of [17]; (3) $S_3(x, \cdot)$ on $\mathbb{R}_+$, for fixed x, uniquely determine v on $\mathbb{R}$ via formulas (2.5), (2.6) and results of [17].*

3   Proofs of Theorems 2.1, 2.2 and 2.3

Proof of Theorem 2.1.
     Using (1.4) we obtain that

$$|\psi^+(x,k)|^2 = \psi^+(x,k)\overline{\psi^+(x,k)} = \tag{3.1}$$

$$1 + 2Re\left(s_{21}(k)e^{-2ikx}\right) + |s_{21}(k)|^2, \quad x \in \mathbb{R}_-, \quad k \in \mathbb{R}_+.$$

In addition, in view of (2.1) we have that

$$2Re\left(s_{21}(k)e^{-2ikx}\right) = 2|s_{21}(k)|\cos(2kx - \alpha(k)). \tag{3.2}$$

Using (2.2), (3.1), (3.2) we obtain that

$$|s_{21}(k)|\cos(2kx)\cos(\alpha(k)) + \sin(2kx)\sin(\alpha(k)) = \tag{3.3}$$

$$2^{-1}(a(x,k) - |s_{21}(k)|^2), \quad x \in \mathbb{R}_-, \quad k \in \mathbb{R}_+.$$

Using (3.3) for $x = x_1$ and $x = x_2$, we obtain the system

$$\begin{pmatrix} \cos(2kx_1) & \sin(2kx_1) \\ \cos(2kx_2) & \sin(2kx_2) \end{pmatrix} |s_{21}| \begin{pmatrix} \cos\alpha \\ \sin\alpha \end{pmatrix} = \\ 2^{-1}\begin{pmatrix} a(x_1,k) - |s_{21}|^2 \\ a(x_2,k) - |s_{21}|^2 \end{pmatrix}, \tag{3.4}$$

where $|s_{21}| = |s_{21}(k)|$, $\alpha = \alpha(k)$.
     Formula (2.3) follows from (3.4).
     Theorem 2.1 is proved.

     Proof of Theorem 2.2. Subtracting equality (3.3) for $x = x_1$ from equality (3.3) for $x = x_2$ and from equality (3.3) for $x = x_3$, we obtain the system

$$\begin{pmatrix} \cos(2kx_2) - \cos(2kx_1) & \sin(2kx_2) - \sin(2kx_1) \\ \cos(2kx_3) - \cos(2kx_1) & \sin(2kx_3) - \sin(2kx_1) \end{pmatrix} |s_{21}| \begin{pmatrix} \cos\alpha \\ \sin\alpha \end{pmatrix} = \\ 2^{-1}\begin{pmatrix} a(x_2,k) - a(x_1,k) \\ a(x_3,k) - a(x_1,k) \end{pmatrix}, \tag{3.5}$$

where $|s_{21}| = |s_{21}(k)|$, $\alpha = \alpha(k)$.
     One can see that

$$\Delta = \sin(2k(x_3 - x_2)) + \sin(2k(x_2 - x_1)) + \sin(2k(x_1 - x_3)), \tag{3.6}$$

where Δ is the determinant of the system (3.5). In addition, using the formulas

$$\sin\varphi_1 + \sin\varphi_2 = 2\cos\left(\frac{\varphi_1 - \varphi_2}{2}\right)\sin\left(\frac{\varphi_1 + \varphi_2}{2}\right), \tag{3.7}$$

$$\sin(\varphi_1 + \varphi_2) = 2\cos\left(\frac{\varphi_1 + \varphi_2}{2}\right)\sin\left(\frac{\varphi_1 + \varphi_2}{2}\right),$$

$$\sin\varphi_1 + \sin\varphi_2 - \sin(\varphi_1 + \varphi_2) = 4\sin\left(\frac{\varphi_1 + \varphi_2}{2}\right)\sin\left(\frac{\varphi_1}{2}\right)\sin\left(\frac{\varphi_2}{2}\right)$$

for $\varphi_1 = 2k(x_2 - x_1)$, $\varphi_2 = 2k(x_1 - x_3)$, we obtain

$$\Delta = 4\sin(k(x_2 - x_3))\sin(k(x_2 - x_1))\sin(k(x_1 - x_3)). \tag{3.8}$$

Formula (2.4) follows from (3.5), (3.8).
Theorem 2.2 is proved.

Proof of Theorem 2.3. Using (3.1) we obtain that

$$(Re\,(s_{21}(k)e^{-2ikx}) + 1)^2 + (Im\,(s_{21}(k)e^{-2ikx}))^2 = |\psi^+(x,k)|^2. \tag{3.9}$$

Formula (2.5), where $(\cdot)^{1/2} > 0$, follows from (3.9) and the property that $|s_{21}(k)| < 1$, see (1.5).

Using (2.1), (3.1), (3.2) we obtain that

$$\frac{d|\psi^+(x,k)|^2}{dx} = 4k|s_{21}(k)|\sin\,(\alpha(k) - 2kx) = 4kIm\,(s_{21}(k)e^{-2ikx}). \tag{3.10}$$

Formula (2.6) follows from (3.10).
Theorem 2.3 is proved.

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